

Energetic emissions from deconfinement in compact stars and their relation to the critical end point in the QCD phase diagram

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Abstract In this work we study the case of deconfinement in compact star interiors in the presence of a strong first order phase transition associated to a critical end point in the QCD phase diagram. Neutron stars fulfilling these conditions show a third branch in the mass-radius diagram with the first and second branches being the white dwarfs and neutron stars configurations. The transition to the third branch can be reached by a pure hadronic neutron star through an induced collapse releasing energy that corresponds to a mass-energy difference between the second and third branch configurations. Physical outcomes of this phenomenon that can potentially explain the already detected astrophysical signals are discussed. In particular we present energy estimations for the case of a fast radio burst, seen as a double-peak structure in the object's light curve.

Keywords: Compact Stars, Deconfinement, Fast Radio Bursts, Massive Twins, Energy Bursts

1. Introduction

Neutron stars (NS) are evolved stars being created after the death of a massive star via a supernova explosion or a transition from a white dwarf accretion or dynamical instabilities. Their interiors can reach up to several times the saturation density n_0 , the canonical density inside atomic nuclei. It is quite uncertain what is the state of matter under such high density conditions, therefore research on equation of state (EoS) is currently a very active area. It is important to note that in NS temperature doesn't play a major role in the computation of the EoS, thus it can be neglected. Recent observations have completely changed our understanding of the cold, dense nuclear matter in such compact star interiors.

In this context, accurate mass determination has proved to be of great importance. In particular the observation of the $2M_{\odot}$ pulsars, PSR J0348+0432 [1] and PSR J1614-2230 [2], have strongly constrained the stiffness of the NS EoS. On the contrary, radius measurements are not yet precise enough to test, discard, and select some of the many alternative EoS models. For instance, frequency resolved pulse shape analysis for the nearest millisecond pulsar [3] supports relatively large radii while analysis from X-Ray bursters [4] point out to either moderate or small radii. Promising future radius measurements include upcoming space missions, cf. NICER [5].

Energetic phenomena like fast radio bursts (FRB) – millisecond duration radio bursts from cosmological distances [6] – can be explained by collapse of compact stars into black holes [7]. Under such scheme, the released energy of the process is about 10^{40} erg. This scenario is not to be confused by early NS evolution consisting of hot protonneutron stars, with deconfinement transition in their interior that could serve as the mechanism triggering the core-collapse supernovae [8], as well as pulsar kicks at birth [9, 10, 11, 12, 13] associated to more energetic emissions.

Neutron stars can suffer a dynamical collapse caused by a deconfinement phase transition in their

cores (possibly leading to a corequake) via spin-down due to continuous electromagnetic emission or by mass accretion. Some early works estimated the energy reservoir for the typical ($\sim 1.4M_{\odot}$) NS mass [14, 15, 16, 17], however here we present a recent EoS model derived in [18] that allows for the formation of a “third family” of compact stars near the maximum mass. Figure 1 shows both gravitational and baryonic masses for compact star configurations with a transition at high masses. We denote the two stars before and after the transition as “high mass twins” because we assume that they bear the same baryonic mass while differing in their mass-energy quantity (binding energy). Although the detailed mechanism is at the moment being developed [19], we can conjecture the following: the deconfinement phase transition occurs via a corequake scenario in which a high-mass hadronic NS collapses into a hybrid compact star disconnected from the former by a gap in the stable configurations (for a recent classification of hybrid stars, see [20]). The instability sets in by the pure hadronic NS after dipole-emission spin-down, or accretion-induced spin-up by matter from a companion.

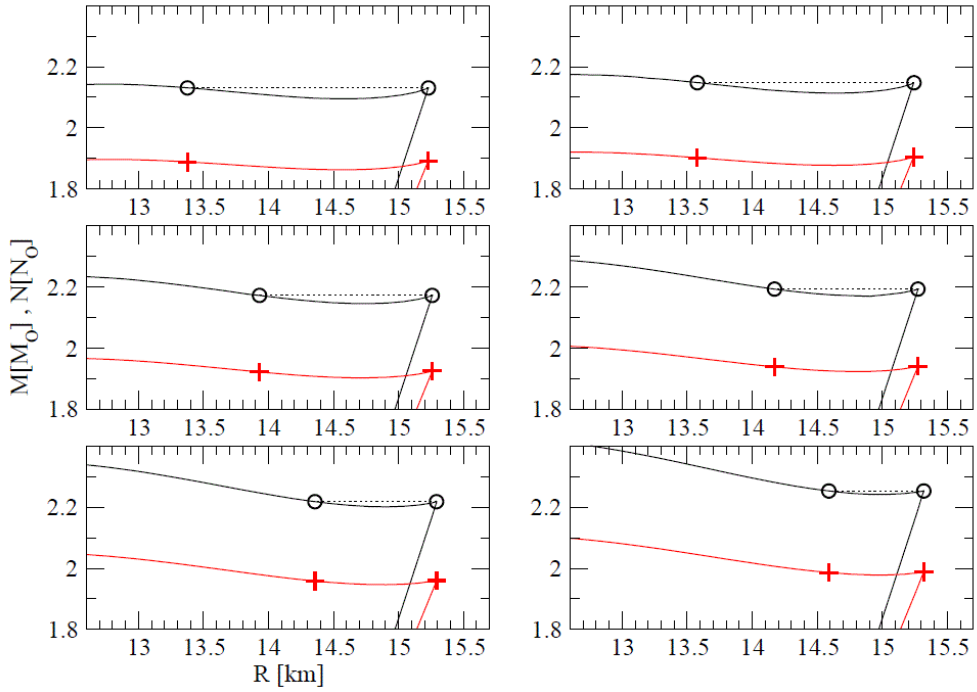


Fig1. Compact star configurations following the instability due to a strong first-order phase transition. Baryonic mass and gravitational mass are given by the black and red curves, respectively. The models are characterized by values of the vector coupling parameter η_4 : 0.0, 1.0, 3.0, 5.0, 7.0, 10.0 starting from the upper left corner down to the bottom right. Notice that the higher the η_4 value the higher the mass at the instability.

For these energy bursts to occur it is important to note that in cold neutron stars the temperatures in their interiors are well below the neutrino opacity temperature ($T_{\nu} \sim 1\text{MeV}$) such that the free streaming scenario applies for neutrino propagation as opposed to the neutrino diffusion mostly suitable for a GRB scenario involving a hot (proto) neutron star ($T_{\nu} < 1\text{MeV}$). A new aspect of this mechanism, going beyond the scenario of Falcke & Rezzolla (a direct collapse of a magnetized NS to a black hole [7]) is the existence of a metastable state between the initial and final states of a supermassive rotating neutron star SURON that could explain a double peak structure of the FRB’s for which recently a case has been made [21]. This metastate in the SURON process corresponds to an object on the third family branch of high

mass hybrid stars, as found recently for microscopically motivated EoS [7, 22].

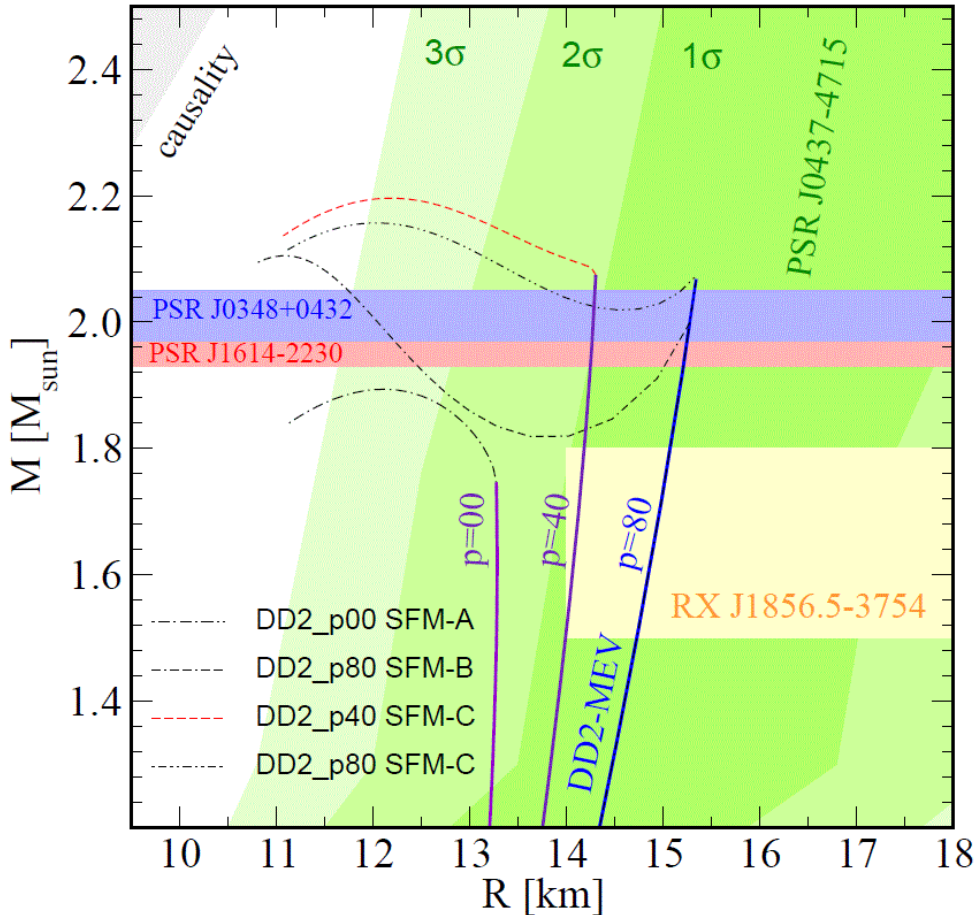


Fig2. String-Flip approach to the high-mass twins [22] and constraints for mass and radius values. The green regions correspond to 1, 2 and 3 σ values as reported in [3]. The vertical bands around $2M_{\odot}$ correspond to the most precisely massive NS measurements [1, 2]. The figure is taken from [22].

2. High-mass twin equation of state

We call neutron star twins those stars with the same mass having different composition and thus different radius. They belong to different branches in the disconnected mass-radius diagram. One of them is pure hadronic whereas the twin star is a hybrid containing quark matter in its core [23, 24, 25, 26, 27]. The case of high mass twins is of great importance because is a consequence of a critical point in the QCD phase diagram [28, 29]. Furthermore, they do not lead to many of the modern issues of compact star physics as discussed in [30]: masquerades [31], the hyperon puzzle [32], and reconfinement [33, 34]. Bayesian studies based on the most reliable observations provide a useful assessment of high mass twins identification [35, 36, 37, 38]. The description of the EoS used here is as follows:

- For hadronic matter we utilize the DD2 (density-dependent relativistic mean field model) EoS

with excluded volume correction (resulting from the internal compositions of nucleons produced by Pauli blocking effects of quarks [39]) that acts at supersaturation densities and provides stiffening to this EoS. For neutron stars this means that the highest mass can be reached at a rather lower interior density values and are characterized by large radii ($\approx 13 - 15$ km, [29]).

- Quark matter is described by a NJL (Nambu–Jona-Lasinio) EoS with multi-quark interactions hNJL [40] featuring a coupling strength parameter in the vector channel of the 8-quark interaction η_4 bringing this quark EoS to a sufficient stiffening at high densities in order to support the maximum observed star mass of $2M_\odot$. Alternatively, an equation of state based on the string-flip model [22] captures the same features (excluded volume), but has the advantage of resulting in a broader range of radius difference between the twins (in better agreement with the contrasting either large or short radius observations). The mass-radius diagram for the string-flip compact stars is shown in Figure 2.

3. Results and discussion

The EoS of [18] features the first order phase transitions leading to instability of the NS right after the appearance of a small, dense quark core. When an NS reaches the maximum hadronic mass radial oscillations take over creating an instability (see e.g. [41]), which results into a dynamical collapse. The criterion for unstable configurations is $\partial M/\partial n_c < 0$, where n_c is the central density of the corresponding NS. In Figure 1 the configurations in between circles (maximum hadronic mass NS and its hybrid twin) that have a positive slope are thus forbidden. This is the origin of a gap created by such prohibited configurations in the mass-radius diagram. Figure 3 (left panel) shows the radius change ΔR in the transition, which is between 1 and 1.5 km for some coupling constant values. The available energy ΔE following the transition equals the mass-energy difference $\Delta M c^2$ between the initial and final configurations and is of about 10^{51} erg. See Figure 3 (right panel). The more realistic case of rotation configurations are being currently studied [19].

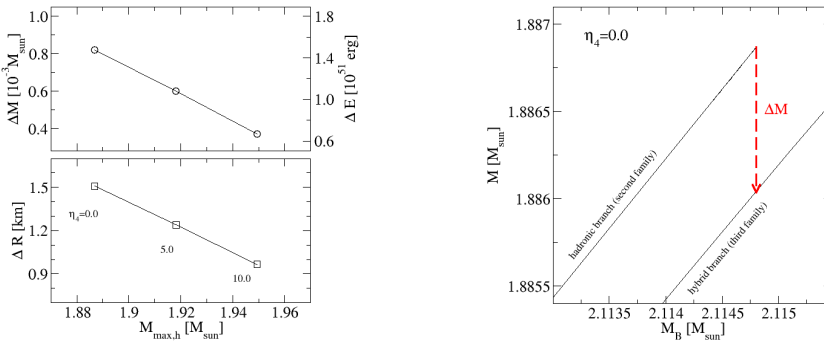


Fig3. Left: Mass difference ΔM (upper panel) and radius difference ΔR (lower panel) resulting from the dynamical NS collapse induced by a deconfinement phase transition for a set of vector coupling parameters $\eta_4 = 0.0, 5.0, 10.0$ of the high-mass twin hNJL models. The corresponding energy release ΔE after the transition is indicated on the right side of the upper panel. Right: Gravitational mass M vs. baryonic mass M_B . $\eta_4 = 0.0$ is the value of the vector coupling in the chosen hNJL model shown here. The red line indicates the energy released in the transition between the maximum hadronic NS and the corresponding hybrid star mass twin. The figures are taken from [42].

We conclude this work by drawing attention to the possibility of energetic emissions produced by a deconfinement phase transition in NS interiors most likely via FRB's, such as the case of the recent

observations of FRB121002 [21] featuring a double peak light curve structure. We identify this double peak signal with the metastable state in the dynamical collapse scenario of a SURON. Furthermore, it is of great importance to mention that the NS EoS presented here can serve as an input to understand scenarios of cosmic ray generation via supernova explosions or NS mergers.

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References

- [1] Antoniadis, J., Freire, P. C. C., Wex, N., et al. 2013, *Science*, 340, 448
- [2] Demorest, P. B., Pennucci, T., Ransom, S. M., Roberts, M. S. E., & Hessels, J. W. T. 2010, *Nature*, 467, 1081
- [3] Bogdanov, S. 2013, *ApJ*, 762, 96
- [4] Steiner, A. W., Lattimer, J. M., & Brown, E. F. 2010, *ApJ*, 722, 33
- [5] Arzoumanian, Z., Gendreau, K. C., Baker, C. L., et al. 2014, *Proc. SPIE*, 9144, 914420
- [6] Lorimer, D. R., Bailes, M., McLaughlin, M. A., Narkevic, D. J., & Crawford, F. 2007, *Science*, 318, 777
- [7] Falcke, H., & Rezzolla, L. 2014, *A&A*, 562, A137
- [8] Sagert, I., Fischer, T., Hempel, M., Pagliara, G., Schaffner-Bielich, J., Mezzacappa, A., Thielemann, F. K. & Liebendorfer, M. 2009, *Phys. Rev. Lett.*, 102, 081101
- [9] Podsiadlowski, P., Pfahl, E., & Rappaport, S. 2005, *Binary Radio Pulsars*, 328, 327
- [10] Lai, D., Chernoff, D. F., & Cordes, J. M. 2001, *ApJ*, 549, 1111
- [11] Kusenko A. 2005, *Int. J. Mod. Phys. A*, 20, 1148 2001, *ApJ*, 549, 1111
- [12] Berdermann, J., Blaschke, D., Grigorian, H., & Voskresensky, D. N. 2006, *Prog. Part. Nucl. Phys.* 57, 334
- [13] Stasielak, J., Biermann, P. L., & Kusenko, A. 2007, *ApJ*, 654, 290
- [14] Drago, A., & Tambini, U. 1999, *J. Phys. G*, 25, 971
- [15] Berezhiani, Z., Bombaci, I., Drago, A., Frontera, F., & Lavagno, A. 2002, *Nucl. Phys. B Proc. Suppl.*, 113, 268
- [16] Aguilera, D. N., Blaschke, D., & Grigorian, H. 2004, *A&A*, 416, 991
- [17] Zdunik, J. L., Bejger, M., Haensel, P., & Gourgoulhon, E., 2008, *A&A*, 479, 515
- [18] Benic, S., Blaschke, D., Alvarez-Castillo, D. E., Fischer, T., & Typel, S. 2015, *A&A*, 577, A40
- [19] Bejger, M., et al., 2016, in preparation.
- [20] Alford, M. G., Han, S., & Prakash, M., 2013, *Phys. Rev. D* 88, 083013

- [21] Champion, D., 2015, talk at seventh Bonn Workshop on “Formation and Evolution of Neutron Stars”, May 18, 2015
- [22] Alvarez-Castillo, D. E., Kaltenborn, M. A. R., & Blaschke, D. 2015, *J. Phys.: Conf. Ser.* 668 (2016) 012035
- [23] Gerlach, U. H. 1968, *Phys. Rev.*, 172, 1325
- [24] Kampfner, B. 1981, *J. Phys. A*, 14, L47
- [25] Schertler, K., Greiner, C., Schaffner-Bielich, J., & Thoma, M. H. 2000, *Nucl. Phys. A*, 677, 463
- [26] Glendenning, N. K., & Kettner, C. 2000, *A&A*, 353, L9
- [27] Dexheimer, V., Negreiros, R., & Schramm, S. 2015, *Phys. Rev. C*, 91, 055808
- [28] Alvarez-Castillo, D. E., & Blaschke, D. 2015, *PoS CPOD 2014* (2014) 045
- [29] Blaschke, D., Alvarez-Castillo, D. E., & Benic, S. 2013, *PoS CPOD 2013* (2013), 063
- [30] Blaschke, D., & Alvarez-Castillo, D. E. 2015, [arXiv:1503.03834](https://arxiv.org/abs/1503.03834)
- [31] Alford, M., Braby, M., Paris, M., & Reddy, S. 2005, *ApJ*, 629, 969
- [32] Baldo, M., Burgio, G. F., & Schulze, H. -J. 2003, *Superdense QCD Matter and Compact Stars*, D. Blaschke and D. Sedrakian (Eds.), Springer, Heidelberg, 2006, p.113
- [33] Lastowiecki, R., Blaschke, D., Grigorian, H., & Typel, S., 2012, *Acta Phys. Polon. Supp.* 5, 535
- [34] Zdunik, J. L., & Haensel, P. 2013, *A&A*, 551, A61
- [35] Ayriyan, A., Alvarez-Castillo, D. E., Blaschke, D., Grigorian, H., & Sokolowski, M. *Particles and Nuclei*, 46 (5) pp. 854-857 (2015)
- [36] Blaschke, D. B., Grigorian, H. A., Alvarez-Castillo, D. E., & Ayriyan, A. S. 2014, *J. Phys. Conf. Ser.*, 496, 012002
- [37] Alvarez-Castillo, D. E., Ayriyan, A., Blaschke, D., & Grigorian, H. 2015, [arXiv:1506.07755](https://arxiv.org/abs/1506.07755). Conference: C14-09-26
- [38] Ayriyan, A., Alvarez-Castillo, D. E., Blaschke, D., & Grigorian, H. 2015, *J. Phys.: Conf. Ser.* 668 (2016) 012038
- [39] Ropke, G., Blaschke, D., & Schulz, H., 1986, *Phys. Rev. D* 34, 3499
- [40] Benic, S. 2014, *Eur. Phys. J. A*, 50, 111
- [41] Shapiro, S. L., Teukolsky, S. A., *Black Holes, White Dwarfs and Neutron Stars: The Physics of Compact Objects*, 1986, Wiley, New York
- [42] Alvarez-Castillo, D. E., Bejger, M., Blaschke, D., Haensel, P., & Zdunik, L. 2015, [arXiv:1506.08645](https://arxiv.org/abs/1506.08645)